How to quickly cool a glass? The Mpemba effect in rugged landscapes

Francesco Ferraro

Laboratory of Interdisciplinary Physics National Biodiversity Future Center University of Padova Italy

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- Z. Lu and O. Raz Proceedings of the National Academy of Sciences 114.20 (2017): 5083-5088
- M. Baity-Jesi et al. The Mpemba effect in spin glasses is a persistent memory effect

Nonequilibrium thermodynamics of the Markovian Mpemba effect and its inverse

Proceedings of the National Academy of Sciences 116.31 (2019): 15350-15355

The Mpemba effect in water



[J. H. Thomas, Wooster Open Works (2007)]



The quasi-static origin of the paradox

distribution space



Non-equilibrium shortcuts



Non-equilibrium Markovian dynamics

Markov property implies master equation

$$\frac{dp_i(t)}{dt} =$$

 Assuming ergodicity + detailed balance: $R_{ii}(T_b)$ is diagonalizable with real eigenvalues

$$\lambda_0 = 0 >$$

and (right) eigenvectors

 $= \sum R_{ij}(T_b)p_j(t)$

 $\lambda_1 > \ldots > \lambda_n$

 $v_0 = p_{eq}, v_1, \dots, v_n$

The Markovian Mpemba effect

- Starting conditions: $p^{h}(0)$ and $p^{c}(0)$
- After quench, relaxation to equilibrium is identical at long times

$$p^{h}(t) = p_{eq} + a_{1}^{h} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} e^{\lambda_{1} t} v_{1} + p_{eq}^{c}(t) = p_{eq}^{c} + a_{1}^{c} + p_{eq}^{c} + p_{eq}^{c} + p_{eq}^{c} + p_{eq}^{c} + p_{eq}^{c} + p_{eq}^{$$

- If $a_1^c > a_1^h$ then the cold system lags behind the hot one: Mpemba effect!
- If $a_1 = 0$, relaxation is "exponentially" faster: "strong" Mpemba effect





Example: three-state system





periodic boundary >θ



Example: three-state system



The first experiment: a colloid in a double-well



[Kumar and Bechhoefer, Nature 2020]

Mpemba effect and rugged landscapes



A simple spin-glass

$$H = -\sum_{\langle i,j \rangle} J_{ij}s$$

- In 3D, spin-glass transition at $T_c = 1.102(3)$
- Simulation details
 - 3D with L = 160
 - 16 samples of $\{J\}$, 256 replicas each
 - Metropolis algorithm, 1 lattice sweep \approx 1 ps
 - simulation time from 1 ps to 0.1 s (!!)

$S_i S_j \qquad J_{ii} = \pm 1$

A realization at $T < T_c$





 $\left\{q_{\boldsymbol{x}}^{(a,b)} = s_{\boldsymbol{x}}^{(a)} \cdot s_{\boldsymbol{x}}^{(b)}\right\}$

Relaxation after a quench below T_c



Relaxation after a quench below T_c



Coherence length

 $s^{(a)}_{oldsymbol{x}}$ $\left\{q_{\boldsymbol{x}}^{(a,b)} = s_{\boldsymbol{x}}^{(a)} \cdot s_{\boldsymbol{x}}^{(b)}\right\}$ $s^{(b)}_{oldsymbol{x}}$



- Above T_c grows (slowly) to equilibrium value
- Below T_c grows without bounds (slowly) as $\xi \sim t^{1/z(T)}$

How to quickly cool a glass?

$E(t) = E_{\infty}(T) + \frac{E_1}{\xi(t)^{D_l}} + \dots$

$D_l \approx 2.5$

lower critical dimension



What about ecological communities?

Editors' Suggestion

$$\frac{dN_i}{dt} = N_i \left[1 - N_i - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \eta_i(t)$$

 $\langle \eta_i(t)\eta_j(t')\rangle = 2TN_i(t)\delta_{ij}\delta(t-t')$

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Properties of Equilibria and Glassy Phases of the Random Lotka-Volterra Model with Demographic Noise

